Top-pair production at hadron colliders

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Top-pair production ... Alexander Mitov Fermilab, 21 Jan 2009

Precision in the LHC era

Precision = confidence!

- LO (leading order) = crude estimate of the result
- NLO (next to leading order) = better estimate of the result crude estimate of uncertainty
- ♦ NNLO = for the first time quantify the uncertainty

Three precision observables have been identified for the LHC:

"The three pillars"

Precision in the LHC era: Drell-Yan

Drell-Yan lepton pair and vector boson production.

✓ luminosity measurement

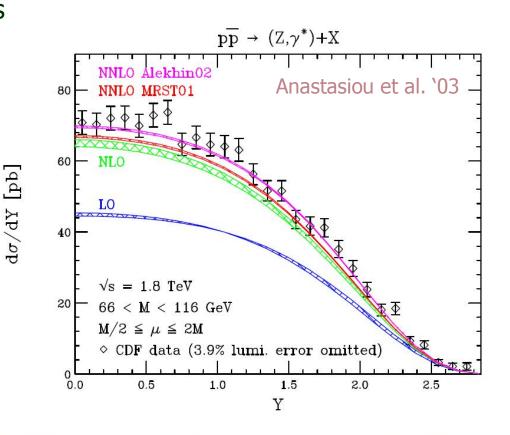
DY is a SM process!

✓ determination of PDF's

NNLO rapidity distribution at the Tevatron.

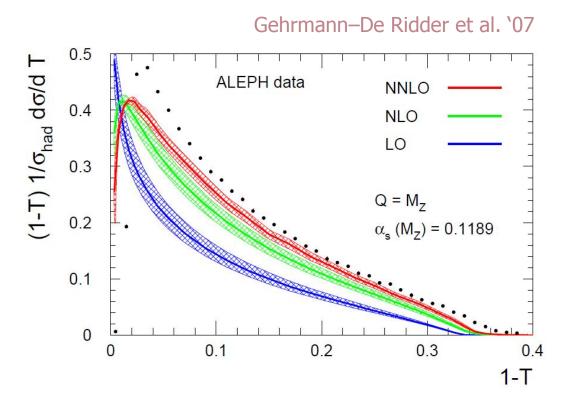
Scale variation:

$$M/2 < \mu < 2M$$



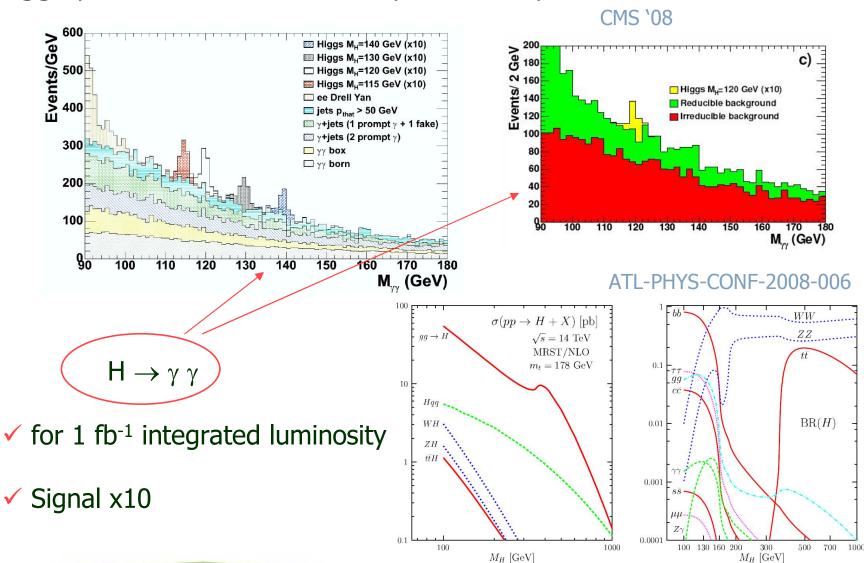
Blast from the past: LEP

Thrust distribution in e⁺e⁻ \rightarrow 3 jets at NNLO Scale variation $M_Z/2 < \mu < 2M_Z$



Precision in the LHC era: Higgs

Higgs production. Its discovery is no simple matter!



Top is central for just about everything:

We want to study it directly (SM)

but also

- Very important for reducing uncertainties in other observables
 - ✓ constrain gluon PDF at larger x,
 - ✓ reduce uncertainties in anti-correlated observables: Examples: singe top and heavy Higgs

CTEQ '08

- Important for Higgs
 - \diamond direct background (in H \rightarrow W+W⁻),
 - conceptually large Yukawa coupling!

Continues ...

- Central collider signature for BSM physics:
 - \star $t\bar{t}$ is a preferred decay mode of TeV KK particles

Top mass (direct/x-section)

Alternative determination of M_{top} from cross-section.

→ Great, because we would be confident in *what* we measure!

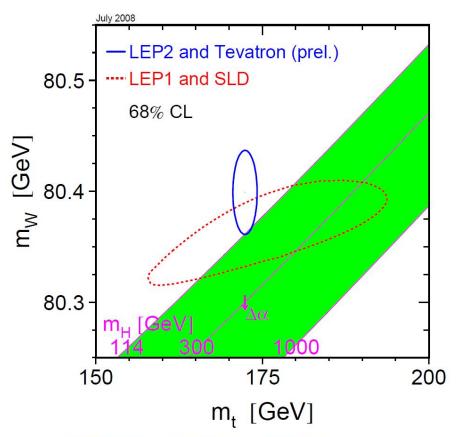
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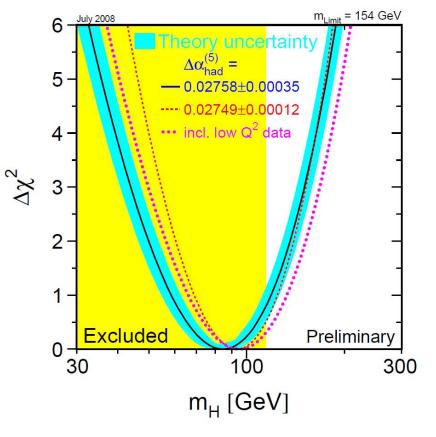
Places where the top mass is crucial: - Higgs mass

- Higgs-inflation

$$M_{\rm t} = 172.4 \pm 1.2 \text{ GeV/}c^2 \text{ CDF+D0 '08}$$

Precision Electroweak Measurements and Constraints on the Standard Model arXiv:0811.4682v1 [hep-ex]





- Top is the most challenging collider signal:
 - 2-to-2 reaction (unlike Drell-Yan and Higgs),
 - involves four colored particles: color correlations
 - massive: its mass cannot be approximated as very small or very large
- The above have led to many theoretical developments
 - prediction of massive amplitudes at high energy,
 - exponentiation, resummation and singularities of massive amplitudes.
- These results can be applied to many more places (jets, tt-Higgs, massless amplitudes at higher orders, N=4 SYM etc.)

An incredible, real life example for

why precision and detailed understanding are very important

The history of the top-pair cross-section calculation:

- ❖ Calculated at NLO ~ 20 years ago (numeric approx.)
- ❖ NLL resummation was done ~ 10 years ago
- It seemed: "top was a done deal";
- It was even called "a plain vanilla".

But let's have a closer look:

The previous results were based on numerical fits (high quality 1%)

Nason, Dawson, Ellis '88 Beenakker et al '89

Only now the exact result could be calculated (2 months ago)

M. Czakon, AM '08

First exact result beyond LO in a 2-to-2 reaction with masses

Compare the new analytic result with the earlier numerical ones.

Of special interest is the threshold expansion $\beta \to 0$ (i.e. $4m^2 \to s)$:

$$f_{gg}^{(1)}(\beta) = \frac{1}{4\pi^{2}} f_{gg}^{(0)}(\beta) \left(\left(C_{F} - \frac{\left(N^{2} - 4 \right) C_{A}}{2(N^{2} - 2)} \right) \frac{\pi^{2}}{2\beta} + 2C_{A} \log^{2}\left(8\beta^{2}\right) - \frac{(9N^{2} - 20)C_{A}}{N^{2} - 2} \log\left(8\beta^{2}\right) + C_{A} \left(\frac{21N^{2} - 50}{N^{2} - 2} - \frac{\left(17N^{2} - 40\right)\pi^{2}}{24(N^{2} - 2)} + \frac{\left(N^{2} - 4\right)\log2}{N^{2} - 2} - 2\log^{2}2 \right) + C_{F} \left(-5 + \frac{\pi^{2}}{4} \right) + O(\beta) \right).$$

It is possible to predict (through *resummation*) the LL,NLL,... behavior of the cross-section close to threshold. Helpful in establishing the theoretical uncertainties.

Extraction of the constant C_3 in the threshold limit:

It is so important, there is an established notation for it

Output

Description:

$$C_A \left(\frac{21N^2 - 50}{N^2 - 2} - \frac{\left(17N^2 - 40\right)\pi^2}{24\left(N^2 - 2\right)} + \frac{\left(N^2 - 4\right)\log 2}{N^2 - 2} - 2\log^2 2 \right) + C_F \left(-5 + \frac{\pi^2}{4}\right)$$

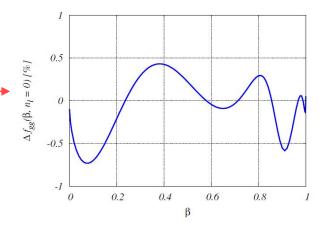
$$= \frac{1111}{21} - \frac{283\pi^2}{168} + \frac{15\log 2}{7} - 6\log^2 2 \approx 34.88,$$

$$\frac{768\pi}{7}a_0^{gg} \approx 37.25$$
.

Czakon, AM '08

Nason, Dawson, Ellis '89

- > X-section better than 1%
- \triangleright But the constant C_3 is 7% different!
- Turns out, it is all consistent ...



Very significant (and unexpected) effect for threshold resummation!

From resummation, the following 2 loop logs can be predicted:

$$\sigma_{gg}(\beta) = \sigma_{gg}^{\text{Born}}(\beta) + \frac{\alpha_s}{4\pi}\sigma_{gg}^{(1)} + \left(\frac{\alpha_s}{4\pi}\right)^2\sigma_{gg}^{(2)} + \mathcal{O}(\alpha_s^3)$$

$$\sigma_{gg}^{(2)} = \sigma_{gg}^{Born}(\beta) \left(4608 \log^4 \beta + 1894.9 \log^3 \beta - 3.4811 \log^2 \beta + \mathcal{O}(\log \beta) \right)$$

Moch Uwer '08

It turns out the coefficient of $\ln^2(\beta)$ is of the form: $(-14306.9505 + 384C_3)$

$$-14306.9505 + 384C_3$$

where: $C_3 = 37.23$

As extracted from NDE '89 and used in ALL resummation literature

 $C_3 = 34.88$

The exact value just recently derived Czakon, AM '08

Therefore the coefficient of $ln^2(\beta)$ becomes -912.35

Note: the reason is pure numerics!

i.e. a change by a factor of 260!

This is a 5% modification to σ , comparable with the current conservative estimates (Cacciari et al '08) of the uncertainty and rather large when compared to (Moch, Uwer '08).

The changes discussed so far are purely due to numerics.

However: there is another modification compared to earlier literature

Exponentiation in Mellin space: $f(N) = \int_0^1 \rho^{N-1} f(\rho) d\rho$, $\rho = 4m^2/s$

$$\sigma_{ij}^{\text{TOT}}(N) = \sigma_{ij,\mathbf{I}}(N) + \sigma_{ij,\mathbf{8}}(N)$$

$$\sigma_{ij,\mathbf{I}}(N) = \sigma_{ij,\mathbf{I}}^{\text{Born}}(N) \ \sigma_{ij,\mathbf{I}}^{\text{H}} \Delta_{ij,\mathbf{I}}(N)$$

We were the first to point out σ^H depend on the color state of the heavy quark pair. We calculated the two coefficients.

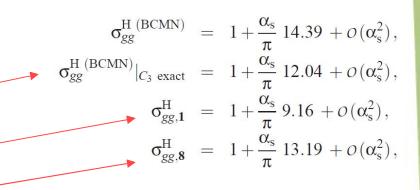
Change in the gg Sudakov resummed

X-section: compare to Bonciani et al '98

C₃ numerics: -5%,

color singlet channel: -12%,

color octet channel: -3%,



Conclusions

- Great time to work in high energy physics!
- LHC will revolutionize our knowledge of the world
 - Discover the Higgs, or its relatives ©
 - Discover New Physics: SUSY, extra dimensions
 - Improve understanding about dark matter and early Universe.
- Successful LHC program requires precision
 - to be confident in what we measure and how we interpret it.
- The challenges ahead are great, and we are more motivated than ever!

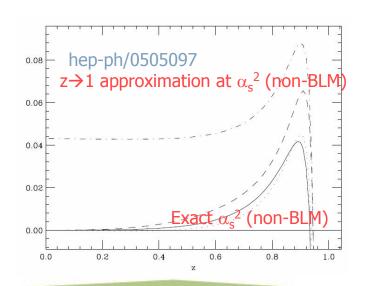
Stay tuned!

What one *needs* to know about top production?

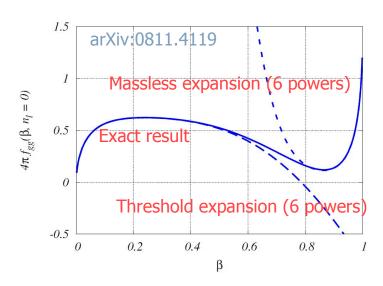
- Understanding true scale uncertainty requires full NNLO calculation !
- The appropriate observable is the total inclusive cross-section.
- Some NNLO terms can be obtained by truncating all-order resummation.
 is this a systematic approximation?

In general, this is a poor approximation to fixed order calculations:

Photon spectrum in $B \rightarrow s + \gamma$:



Top X-section: NLO correction



Back to Phenomenology

What our results actually imply?

Conclusion #1: the earlier FO NLO calculations are of high quality

However: in the last 10 years or so, all "improvements" have been done by soft gluon resummation (since no new FO results)

Idea of threshold resummation: towers of $ln(\beta)$ terms in the total inclusive cross section $\sigma_{TOT}(\beta)$ can be predicted to all orders in the coupling.

Conclusion #2: the resummation results are altered by our work